

1. Pembuktian Rumus

$$V_{coulomb} = \frac{-\alpha e^2}{4\pi\epsilon_0 r} \text{ secara pendekatan : } V_{rep} = \frac{B}{r^n}$$

Sehingga energy potensial total :

$$V = V_{coulomb} + V_{rep}$$

$$V = \frac{-\alpha e^2}{4\pi\epsilon_0 r} + \frac{B}{r^n} \dots\dots\dots (1)$$

Jarak kesetimbangan ion-ion terjadi bila V mempunyai harga minimum

V_0 pada $r = r_0$

$$\frac{dV}{dr} = 0$$

$$\frac{d}{dr} \left(\frac{-\alpha e^2}{4\pi\epsilon_0 r} + \frac{B}{r^n} \right) = 0$$

$$\frac{-\alpha e^2}{4\pi\epsilon_0 r} (-1)r^{-2} + B(-n)r^{(-n-1)} = 0, \quad r = r_0$$

$$\frac{-\alpha e^2}{4\pi\epsilon_0 r_0^2} - \frac{B}{r_0^{n+1}} = 0$$

$$B = \frac{\alpha e^2}{4\pi\epsilon_0 r_0^2} \frac{r_0^{n+1}}{n} = \frac{\alpha e^2 r_0^{n-1}}{4\pi\epsilon_0 n} \dots\dots\dots (2)$$

Substitusi persamaan 2 ke 1

$$\begin{aligned} V_0 &= \frac{-\alpha e^2}{4\pi\epsilon_0 r} + \frac{1}{r^n} \left(\frac{\alpha e^2}{4\pi\epsilon_0 n} r^{n-1} \right) \\ &= \frac{-\alpha e^2}{4\pi\epsilon_0 r} + \frac{\alpha e^2}{4\pi\epsilon_0 r n} = \frac{\alpha e^2}{4\pi\epsilon_0 r} \left(1 - \frac{1}{n} \right) \end{aligned}$$

2. Kapasitas Panas Model Einstein

$$f(\nu) = \frac{1}{e^{h\nu/kT} - 1}$$

Energi rata-rata nya:

$$\bar{E} = h\nu \cdot f(\nu)$$

$$\bar{E} = \frac{h\nu}{e^{h\nu/kT} - 1}$$

Untuk 1 kilo mol maka energy dalamnya:

$$U = 3N_0\bar{E}$$

$$U = 3N_0 \frac{h\nu}{e^{h\nu/kT} - 1}$$

↳ Untuk temperature tinggi $T \gg \rightarrow \frac{h\nu}{kT} \ll 1$

$$\text{Ingat : } e^x = 1 + x + \frac{x^2}{2!} + \frac{x^3}{3!} + \dots$$

$$e^{h\nu/kT} = 1 + \frac{h\nu}{kT} + \dots$$

Maka :

$$U = 3N_0 \frac{h\nu}{e^{h\nu/kT} - 1} = 3N_0 \frac{h\nu}{\left(1 + \frac{h\nu}{kT}\right) - 1} = \frac{3N_0 h\nu}{\frac{h\nu}{kT}}$$

$$U = 3N_0 kT \rightarrow \text{dimana } N_0 k = R$$

$$U = 3RT$$

$$Cv = \frac{\partial u}{\partial T} = \frac{\partial}{\partial T} (3RT)$$

Maka : $Cv = 3R \rightarrow$ kapasitas panas model Einstein untuk $T \gg$

(sesuai dengan eksperimen Dulong dan Petit)

↳ Untuk suhu rendah $T \ll \rightarrow \frac{h\nu}{kT} \gg 1$

$$U = 3N_0 \frac{h\nu}{e^{h\nu/kT} - 1}$$

$$Cv = \frac{\partial U}{\partial T} = \frac{\partial}{\partial T} \left(3N_0 \frac{h\nu}{e^{h\nu/kT} - 1} \right)$$

$$= 3N_0 h\nu \left(-1 \left(e^{h\nu/kT} - 1 \right)^{-2} \right) \left(\frac{-h\nu}{kT} e^{h\nu/kT} \right)$$

$$Cv = 3N_0 h\nu \frac{h\nu}{kT^2} \frac{e^{h\nu/kT}}{\left(e^{h\nu/kT} - 1 \right)^2}$$

$$Cv = 3N_0 k \left(\frac{h\nu}{kT} \right)^2 \frac{e^{h\nu/kT}}{\left(e^{h\nu/kT} - 1 \right)^2}$$

$$Cv = 3R \left(\frac{h\nu}{kT} \right)^2 \frac{e^{h\nu/kT}}{\left(e^{h\nu/kT} - 1 \right)^2}$$

Temperatur khusus : suhu Einstein $\theta_E = \frac{h\nu}{k}$

Maka : $Cv = 3R \left(\frac{\theta_E}{T} \right)^2 \frac{e^{\theta_E/T}}{\left(e^{\theta_E/T} - 1 \right)^2} \rightarrow$ kapasitas panas model Einstein untuk $T \ll$

Kapasitas panas Model Debye

Rapat keadaan

Di definisikan : $D(W) = \frac{\text{Jumlah keadaan } (dN)}{\text{Rentang energi } (dW)}$

Maka jumlah keadaan : $dN = D(W) dW$

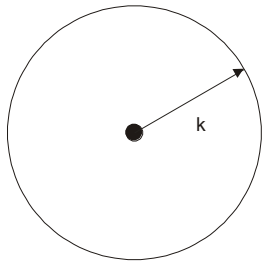
Energi Total

$$U = \sum_k \left\{ \sum_p \frac{\hbar \omega_{kp}}{e^{\frac{\hbar \omega_{kp}}{k_B T}} - 1} \right\}$$

Debye menganggap energinya bersifat continue, maka :

$$U = \sum_p \int \frac{\hbar \omega_{kp}}{e^{\frac{\hbar \omega_{kp}}{k_B T}} - 1} D(W) dW \dots (1)$$

Volume Ruang



$$N = \frac{\frac{4}{3}\pi k^3}{\left(\frac{2\pi}{L}\right)^3}$$

$$N = \frac{4}{3}\pi k^3 \cdot \frac{L^3}{8\pi^3} = \frac{L^3 k^3}{6\pi^2} = \frac{V k^3}{6\pi^2}$$

$$D(k) = \frac{dN}{dk} = \frac{d}{dk} \left(\frac{V k^3}{6\pi^2} \right) = \frac{3V k^2}{6\pi^2}$$

$$D(k) = \frac{dN}{dk} = \frac{V k^2}{2\pi^2}$$

$$D(\omega) = \frac{dN}{d\omega} = \frac{dN}{dk} \cdot \frac{dk}{d\omega} = \frac{V k^2}{2\pi^2} \left(\frac{dk}{d\omega} \right)$$

$$D(\omega) = \frac{V k^2}{2\pi^2} \frac{1}{V} = \frac{V \left(\frac{\omega^2}{V} \right)}{2\pi^2} \frac{1}{V}$$

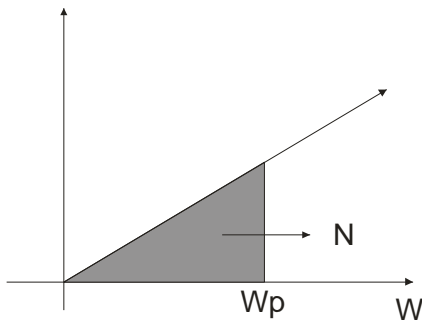
$$D(\omega) = \frac{V \omega^2}{2\pi^2 V} \dots (2)$$

Substitusi persamaan (2) ke (1) :

$$U = \sum_p \int \frac{\hbar \omega_{kp}}{e^{\frac{\hbar \omega_{kp}}{k_B T}} - 1} D(\omega) d\omega$$

$$U = 3 \int_0^{\omega_0} \frac{\hbar \omega}{e^{\hbar \omega / k_B T} - 1} \cdot \frac{V \omega^2}{2\pi^2 V} d\omega$$

$$U = \int_0^{\omega_0} \frac{\hbar \omega^3 V / 2\pi^2 V^3}{e^{\hbar \omega / k_B T} - 1} d\omega$$



$$\omega_D = \omega_{Debye}$$

$$\omega_D = V k_D$$

$$N(\omega) = N_{total}$$

$$N = \frac{\frac{4}{3} \pi k_D^3}{\left(\frac{2\pi}{L}\right)^3} = \frac{L^3 k_D^3}{6\pi^2} = \frac{V k_D^3}{6\pi^2}$$

$$V = \frac{6\pi^2}{k_D^3} N \dots\dots (3)$$

Kapasitas Panas

$$C_v = \frac{\partial U}{\partial T} = \frac{\partial}{\partial T} \left[3 \int_0^{\omega_0} \frac{\hbar \omega}{e^{\hbar \omega / k_B T} - 1} \cdot \frac{V \omega^2}{2\pi^2 V} d\omega \right]$$

$$\begin{aligned}
Cv &= \frac{3\hbar V}{2\pi^2 V^3} \int_0^{w_D} w^3 \left(\frac{d}{dT} \frac{1}{e^{\hbar w/k_B T} - 1} \right) dw \\
&= -(e^{\hbar w/k_B T} - 1)^{-2} \cdot \frac{\hbar w}{k_B T^2} e^{\hbar w/k_B T} \\
&= \frac{\hbar w}{k_B T^2} \frac{e^{\hbar w/k_B T}}{(e^{\hbar w/k_B T} - 1)^2}
\end{aligned}$$

Maka :

$$Cv = \frac{3\hbar V}{2\pi^2 V^3} \int_0^{w_D} w^3 \frac{\hbar w}{k_B T^2} \frac{e^{\hbar w/k_B T}}{(e^{\hbar w/k_B T} - 1)^2} dw$$

$$Cv = \frac{3\hbar^2 V}{2\pi^2 V^3} \frac{1}{k_B T^2} \int_0^{w_D} w^4 \frac{e^{\hbar w/k_B T}}{(e^{\hbar w/k_B T} - 1)^2} dw$$

Misal : $x = \frac{\hbar w}{k_B T} \rightarrow w = \frac{k_B T}{\hbar} x$

$$dw = \frac{k_B T}{\hbar} dx$$

$$Cv = \frac{3\hbar^2 V}{2\pi^2 V^3} \frac{1}{k_B T^2} \int_0^x w^4 \frac{\left(\frac{k_B T}{\hbar} x\right)^4 e^x}{(e^x - 1)^2} \frac{k_B T}{\hbar} dx$$

$$Cv = \frac{3\hbar^2 V}{2\pi^2 V^3} \frac{1}{k_B T^2} \int_0^{\hbar w/k_B T} w^4 \frac{\left(\frac{k_B T}{\hbar}\right)^4 x^4 e^x}{(e^x - 1)^2} \frac{k_B T}{\hbar} dx$$

$$= \frac{3\hbar^2 V}{2\pi^2 V^3} \frac{1}{k_B T^2} \frac{k_B T}{\hbar} \left(\frac{k_B T}{\hbar}\right)^4 \int_0^{\hbar w/k_B T} \frac{x^4 e^x}{(e^x - 1)^2} dx$$

$$Cv = \frac{3V k_B^4 T^3}{2\pi^2 V^3 \hbar^3} \int_0^{\hbar w/k_B T} \frac{x^4 e^x}{(e^x - 1)^2} dx$$

Jika $\theta_D = \frac{\hbar w}{k_B}$ ingat $x = \frac{\hbar w}{k_B}$

Maka $\frac{\hbar w}{k_B} = \frac{\theta_D}{T}$

$$Cv = \frac{3Vk_B^4 T^3}{2\pi^2 V^3 \hbar^3} \int_0^{\theta_D/T} \frac{x^4 e^x}{(e^x - 1)^2} dx$$

Dari persamaan (3), $V = \frac{6\pi^2}{k_D^3} N \rightarrow$ ingat $w_D = v k_D$

$$k_D^3 = \frac{w_D^3}{v^3}$$

$$V = \frac{6\pi^2 N v^3}{w_D^3}$$

$$Cv = \frac{3 \left(\frac{6\pi^2 N v^3}{w_D^3} \right) k_B^4 T^3}{2\pi^2 V^3 \hbar^3} \int_0^{\theta_D/T} \frac{x^4 e^x}{(e^x - 1)^2} dx$$

$$Cv = \frac{9Nk_B^4 T^3}{\hbar^3 w_D^3} \int_0^{\theta_D/T} \frac{x^4 e^x}{(e^x - 1)^2} dx, \text{ ingat } \theta_D = \frac{\hbar w}{k_B} \rightarrow w_D = \frac{\theta_D}{\hbar} k_B$$

$$Cv = \frac{9Nk_B^4 T^3}{\hbar^3 \left(\frac{\theta_D}{\hbar} k_B \right)^3} \int_0^{\theta_D/T} \frac{x^4 e^x}{(e^x - 1)^2} dx, \text{ ingat } x = \frac{\theta_D}{T}$$

∇ Untuk suhu tinggi $T \gg \theta_D \rightarrow x_D \ll 1$

Maka :

$$\frac{x^4 e^x}{(e^x - 1)^2} = \frac{x^4 e^x}{(e^x - 1)(e^x - 1)} = \frac{x^4 e^x}{(e^x - 1)(1 - e^x)}$$

Deret :

$$e^x = 1 + x + \frac{x^2}{2!} + \frac{x^3}{3!} + \dots$$

$$e^{-x} = 1 - x + \frac{x^2}{2!} - \frac{x^3}{3!} + \dots$$

$$(e^x - 1)(e^{-x} - 1) = (e^x - 1)(1 - e^{-x})$$

$$= \left(1 + x + \frac{x^2}{2!} + \frac{x^3}{3!} + \dots\right) \left(1 - \left(x + \frac{x^2}{2!} - \frac{x^3}{3!} + \dots\right)\right)$$

$$= \left(x + \frac{x^2}{2!} + \frac{x^3}{3!}\right) \left(x - \frac{x^2}{2!} + \frac{x^3}{3!}\right)$$

$$= x^2 - \frac{x^3}{2!} + \frac{x^4}{3!} + \frac{x^3}{2!} - \frac{x^4}{2!2!} + \frac{x^5}{2!3!} + \frac{x^4}{3!} - \frac{x^5}{2!3!} + \frac{x^6}{3!3!}$$

$$= x^2 + \frac{x^4}{3!} - \frac{x^4}{2!2!} + \frac{x^4}{3!} + \dots$$

$$= x^2 + \frac{2x^4}{3!} - \frac{x^4}{2!2!} = x^2 + x^4 \left(\frac{2}{3 \cdot 2 \cdot 1} - \frac{1}{2 \cdot 1 \cdot 2 \cdot 1}\right)$$

$$(e^x - 1)(1 - e^{-x}) = 2 \left(\frac{x^2}{2} + \frac{x^4}{24}\right) = 2 \left(\frac{x^2}{2!} + \frac{x^4}{4!} + \dots\right)$$

$$x_D \ll 1 \rightarrow \frac{x^4 e^x}{(e^x - 1)^2} = \frac{x^4 e^0}{(e^x - 1)(1 - e^{-x})} = \frac{x^4}{2 \left\{ \frac{x^2}{2!} + \frac{x^4}{4!} + \dots \right\}} = \frac{x^4}{2 \cdot \frac{x^2}{2!}} \approx x^2$$

Jadi :

$$Cv = \frac{9Nk_B T^3}{\theta_D^3} \int_0^{\theta_D} x^2 dx = 9Nk_B \left(\frac{T}{\theta_D}\right)^3 \frac{1}{3} x^3 \rightarrow \text{ingat } x = \frac{\theta_D}{T}$$

$$Cv = 9Nk_B \left(\frac{T}{\theta_D}\right)^3 \frac{1}{3} \frac{\theta_D^3}{T^3}$$

$$Cv = 3Nk_B \rightarrow \text{dimana } Nk_B = R$$

$$Cv = 3R \rightarrow \text{kapasitas panas model Debye untuk } T \gg$$

↳ Untuk suhu rendah $T \ll \theta_D \rightarrow x_D \gg 1$

$$Cv = 9Nk_B \left(\frac{T}{\theta_D}\right)^3 \int_0^{\theta_D/T} \frac{x^4 e^x}{(e^x - 1)^2} dx$$

↓
①

① → integral parsial $\int u dv = uv - v \int du$

$$\int_0^{\theta_D/T} \frac{x^4 e^x}{(e^x - 1)^2} dx \rightarrow \text{misal } u = x^4$$

$$du = 4x^3 dx$$

$$dv = \frac{e^x}{(e^x - 1)^2} dx \rightarrow v = \int \frac{e^x}{(e^x - 1)^2} dx$$

$$\text{Misal: } e^x - 1 = u$$

$$du = e^x dx$$

$$v = \int \frac{1}{u^2} du = \int u^{-2} du$$

$$v = -\frac{1}{u} = -\frac{1}{(e^x - 1)}$$

$$\int_0^{\theta_D/T} \frac{x^4 e^x}{(e^x - 1)^2} dx = \frac{-x^4}{(e^x - 1)} + 4 \int \frac{x^3}{(e^x - 1)} dx$$

$$\text{jadi : } Cv = 9Nk_B \left(\frac{T}{\theta_D}\right)^3 \left\{ \frac{-x^4}{(e^x - 1)} + 4 \int_0^{\theta_D/T} \frac{x^3}{(e^x - 1)} dx \right\}$$

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②

↓

③

$$\text{②} \rightarrow \text{untuk } T = 0 \text{ maka } \frac{-x^4}{(e^x - 1)} \approx \frac{\left(\frac{\theta_D}{T}\right)^4}{e^{\theta_D/T} - 1} = 0$$

③ → jika $4 \int_0^{\infty} \frac{x^3}{(e^x - 1)} dx \rightarrow$ gunakan $\frac{1}{e^x - 1} = \sum_{\alpha=1}^{\infty} e^{-\alpha x}$

$$4 \int_0^{\infty} \frac{x^3}{(e^x - 1)} dx = 4 \int_0^{\infty} x^3 \sum_{\alpha=1}^{\infty} e^{-\alpha x} dx$$

$$\rightarrow \text{rumus} \int_0^{\infty} x^n e^{-\alpha x} dx = \frac{n!}{\alpha^{n+1}}$$

$$= \left[3! \sum_{\alpha=1}^{\infty} \frac{1}{\alpha^4} dx \right]$$

$$= 4 \left[6 \left(\frac{1}{1^4} + \frac{1}{2^4} + \frac{1}{3^4} + \frac{1}{4^4} + \dots \right) \right]$$

↓

$$\text{Rumus : } \frac{1}{1^{2p}} + \frac{1}{2^{2p}} + \frac{1}{3^{2p}} + \frac{1}{4^{2p}} + \dots = \frac{2^{p-1} \pi^{2p}}{(2p)!}$$

Didapat $p = 2 \rightarrow$ untuk $p = 2 \rightarrow B_p = \frac{1}{30}$ (Dari tabel bernoulli)

$$\frac{1}{1^4} + \frac{1}{2^4} + \frac{1}{3^4} + \frac{1}{4^4} = \frac{2^3 \pi^4}{4!} \frac{1}{30} = \frac{8\pi^4}{4 \cdot 3 \cdot 2 \cdot 1} \frac{1}{30} = \frac{\pi^4}{90}$$

$$4 \int_0^{\infty} \frac{x^3}{(e^x - 1)} dx = 4 \left[6 \cdot \frac{\pi^4}{90} \right] = 8 \left(\frac{\pi^4}{30} \right) = \frac{4\pi^4}{15}$$

Jadi :

$$Cv = 9Nk_B \left(\frac{T}{\theta_D} \right)^3 \left\{ \frac{-x^4}{(e^x - 1)} + 4 \int_0^{\infty} \frac{x^3}{(e^x - 1)} dx \right\}$$

$$Cv = 9Nk_B \left(\frac{T}{\theta_D} \right)^3 \left\{ 0 + \frac{4\pi^4}{15} \right\}$$

$$Cv = 9Nk_B \left(\frac{T}{\theta_D}\right)^3 \frac{4\pi^4}{15}$$

$$Cv = \frac{12}{5} \pi^4 Nk_B \left(\frac{T}{\theta_D}\right)^3$$

$$Cv = 234Nk_B \left(\frac{T}{\theta_D}\right)^3 \rightarrow \text{kapasitas panas model Debye untuk } T \ll$$

5. Turunkan bahwa kapasitas panas untuk elektron adalah :

$$Cv = \frac{\pi^2 T k_B^2 N}{2E_F}$$

Dari mekanika klasik diperoleh bahwa :

$$\text{Energi untuk satu derajat kebebasan } U = \frac{1}{2} k_B T$$

Sehingga untuk partikel dengan 3 derajat kebebasan adalah:

$$U = 3 \frac{1}{2} k_B T = \frac{3}{2} k_B T$$

$$\text{Dan kapasitas panas untuk 1 partikel adalah : } Cv = \frac{dU}{dT} = \frac{3}{2} k_B$$

$$\text{Maka kapasitas panas untuk N buah partikel : } Cv = \frac{3}{2} Nk_B$$

$$\text{Besarnya energi fermi : } E_F = \frac{\hbar^2}{2m} k_p^2$$

$$\text{Bila } \frac{T}{T_F} \approx 0,01 \text{ maka energi ferminya } E_F = k_B \cdot T_F \rightarrow T_F = \frac{E_F}{k_B}$$

∴ Untuk suhu rendah $k_B T \ll E_F$

$$\text{Distribusi fermi diract : } F(E) = \frac{1}{e^{(E-E_F)/k_B T} + 1}$$

$$T = 0$$

$$E < E_F \rightarrow f(E) = 1$$

$$E > E_F \rightarrow f(E) = 0$$

Bila perubahan energinya adalah $U = U(T) - U(0)$

Dimana $U(T)$ = energi setelah elektron pindah dari keadaan dasar

$$U = U(T) - U(0)$$

$$U = \int_0^2 D(E) - F(E)E dE - \int_0^{E_F} D(E)E dE$$

$$\text{Bila } N = \int_0^{E_F} D(E) F(E) dE = \int_0^{\sim} D(E)F(E) dE$$

$$NE_F = \int_0^{\sim} D(E) F(E)E_F dE = \left[\int_0^{E_F} + \int_{E_F}^{\sim} \right] D(E)F(E)E_F dE$$

Maka :

$$U = \int_{E_F}^{\sim} D(E)F(E)(E - E_F) dE + \int_0^{E_F} D(E)(E_F - E)(1 - F(E)) dE$$

Karena integral yang bergantung pada suhu hanya $F(E)$ maka diferensiasinya terhadap T hanya berlaku untuk suku-suku yang mengandung $F(E)$ saja.

$$\begin{aligned}
Cv &= \frac{dU}{dT} \\
&= \frac{d}{dT} \left[\int_{E_F}^{\sim} D(E)F(E)(E - E_F) dE \right. \\
&\quad \left. + \int_0^{E_F} D(E)(E_F - E) dE + \int_0^{E_F} D(E)(E - E_F) F(E) dE \right] \\
&= \frac{d}{dT} \left[\int_0^{E_F} D(E)(E - E_F) F(E) dE + \int_{E_F}^{\sim} D(E)(E - E_F) F(E) dE \right] \\
&= \frac{d}{dT} D(E)F(E)(E - E_F) dE \\
&= k_B \int_0^{\sim} D(E)(E - E_F) \frac{dE(E)}{dT k_B} dE
\end{aligned}$$

$$Cv = k_B \int_0^{\sim} D(E)(E - E_F) \frac{dF(E)}{dT k_B} dE \dots\dots\dots (1)$$

Ingat : $F(E) = \frac{1}{e^{(E-E_F)/k_B T}}$

Maka $\frac{dF(E)}{dT k_B} = \frac{d}{dT} (e^{(E-E_F)/T} + 1)^{-1}$ dimana $\tau = k_B T$

$$\frac{dF(E)}{d\tau} = \frac{E-E_F}{\tau^2} \frac{e^{(E-E_F)/\tau}}{(e^{(E-E_F)/\tau} + 1)^2} \dots\dots\dots (2)$$

Persamaan (1) menjadi :

$$Cv = k_B \int_0^{\sim} D(E)(E - E_F) \frac{dF(E)}{d\tau} dE$$

Untuk $T \ll$ maka persamaan diatas menjadi :

$$Cv = k_B D(E_F) \int_0^{\sim} (E - E_F) \frac{(E - E_F)}{\tau^2} \frac{e^{(E-E_F)/\tau}}{(e^{(E-E_F)/\tau} + 1)^2} dE$$

$$Cv = k_B D(E_F) \int_0^{\sim} (E - E_F) \frac{(E - E_F)}{\tau^2} \frac{e^{(E-E_F)/\tau}}{(e^{(E-E_F)/\tau} + 1)^2} dE$$

Misal : $x = \frac{(E-E_F)}{\tau}$

$E = 0 \rightarrow x = -E_F/\tau$

$E = \sim \rightarrow x = \sim$

$$Cv = k_B D(E_F) \int_{E_F/\tau}^{\sim} x^2 \frac{e^x}{(e^x + 1)^2} \tau dx \rightarrow \text{ingat } \tau = k_B T$$

$$Cv = k_B D(E_F) \int_{E_F/\tau}^{\sim} x^2 \frac{e^x}{(e^x + 1)^2} k_B T dx$$

$$Cv = k_B^2 T D(E_F) \int_{E_F/\tau}^{\sim} x^2 \frac{e^x}{(e^x+1)^2} dx \dots (4)$$

↓

$$\int_{E_F/\tau}^{\sim} x^2 \frac{e^x}{(e^x+1)^2} dx \rightarrow \text{misal } u = x^2$$

$$du = 2x dx$$

$$dv = \frac{e^x}{(e^x+1)^2} dx$$

$$v = -\frac{1}{(e^x+1)}$$

Gunakan integral parsial

$$\int_{E_F/\tau}^{\sim} x^2 \frac{e^x}{(e^x+1)^2} dx = -\frac{x^2}{(e^x+1)} + 2 \int \frac{x}{(e^x+1)} dx \dots (5)$$

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②

$$\textcircled{1} \rightarrow -\frac{x^2}{(e^x+1)} = -\frac{\left(\frac{(E-E_F)}{\tau}\right)^2}{e^{\frac{-(E-E_F)}{\tau}} + 1} \rightarrow \text{dimana } \tau = k_B T \text{ untuk } T = 0$$

$$\text{Maka } \frac{(E-E_F)}{\tau} = 0$$

Sehingga $-\frac{x^2}{(e^x+1)} = 0$

$$\textcircled{2} \rightarrow 2 \int \frac{x}{(e^x+1)} dx \rightarrow \text{gunakan } \frac{1}{(e^x+1)} = \sum_{\tilde{\alpha}=1}^{\sim} e^{\alpha x}$$

$$2 \int \sum_{\tilde{\alpha}=1}^{\sim} e^{\alpha x} dx \text{ rumus } \int_0^{\sim} x^n e^{\alpha x} dx = \frac{n!}{(-\alpha)^{n+1}}$$

$$= 2 \left[1! \sum_{\tilde{\alpha}=1}^{\sim} \frac{1}{-\alpha^2} \right]$$

$$= 2 \left[1 \left(\frac{1}{(-1)^2} + \frac{1}{(-2)^2} + \dots \right) \right]$$

↓

$$\text{Rumus : } \frac{1}{1^{2p}} + \frac{1}{2^{2p}} + \frac{1}{3^{2p}} + \frac{1}{4^{2p}} + \dots = \frac{2^{p-1} \pi^{2p}}{(2p)!}$$

Didapat $p = 1 \rightarrow$ untuk $p = 1 \rightarrow B_p = \frac{1}{6}$ (Dari tabel bernoulli)

$$\frac{1}{1^2} + \frac{1}{2^2} + \dots = \frac{2^{2-1}\pi^2 \frac{1}{6}}{2!} = \frac{\pi^2}{6}$$

Jadi :

$$2 \left[1! \sum_{\alpha=1}^{\infty} \frac{1}{-\alpha^2} \right] = 2 \left[\frac{1}{1^2} + \frac{1}{2^2} + \dots \right] = 2 \left(\frac{\pi^2}{6} \right) = \frac{\pi^2}{3}$$

$$2 \int \frac{x}{(e^x + 1)} dx = \frac{\pi^2}{3}$$

Substitusi ① dan ② ke persamaan (5) maka diperoleh

$$\int_{E_F/\tau}^{\infty} x^2 \frac{e^x}{(e^x+1)^2} dx = 0 + \frac{\pi^2}{3} = \frac{\pi^2}{3} \dots\dots\dots (6)$$

Substitusi persamaan (6) ke (4)

$$Cv = k_B^2 T D(E_F) \int_{E_F/\tau}^{\infty} x^2 \frac{e^x}{(e^x + 1)^2} dx$$

$$Cv = k_B^2 T D(E_F) \frac{\pi^2}{3} \rightarrow \text{rapat keadaan } D(E) = \frac{3N}{2E_F}$$

Maka :

$$Cv = k_B^2 T \frac{3N}{2E_F} \frac{\pi^2}{3}$$

$$Cv = k_B^2 T \frac{N}{2E_F} \pi^2$$

$$Cv = \frac{\pi^2 T k_B^2 N}{2E_F} \leftarrow \text{Terbukti}$$

6. untuk elektron bebas

$$E_k = \frac{\hbar^2}{2m} (k_x^2 + k_y^2 + k_z^2)$$

Dimana : $k_x, k_y, k_z = 0; \pm \frac{2\pi}{L}; \pm \frac{4\pi}{L}; \dots \dots$

Sehingga :

$$\psi_{(+)} = \exp(i \pi x/a) + \exp(-i \pi x/a) = 2 \cos \frac{\pi x}{a}$$

$$\psi_{(-)} = \exp(i \pi x/a) - \exp(-i \pi x/a) = 2i \sin \frac{\pi x}{a}$$

- Misalkan gelombang berdiri pada persamaan di atas

$$\psi^* \psi = |\psi|^2$$

$$\rho = \exp(-ikx) \exp(ikx) = 1$$

$$\rho_{(+)} = |\psi_{(+)}|^2 \cos^2 \frac{\pi x}{a}$$

$$\rho_{(-)} = |\psi_{(-)}|^2 \sin^2 \frac{\pi x}{a}$$

- Besarnya energi Gap

$$V(x) = V \cos \frac{2\pi x}{a}$$

$$E_g = \int_0^1 dx V(x) \left[|\psi_{(+)}|^2 + |\psi_{(-)}|^2 \right]$$

$$= 2 \int dx V \cos(2x\pi/a) (\cos^2 \pi x/a - \sin^2 \pi x/a) = 0$$

$$\psi_{(x+a)} = c \psi_{(x)}$$

$$\psi_{(x+Na)} = \psi_{(x)} = c^N \psi_{(x)}$$

$$c = \exp(i 2\pi s/N); \quad s = 0, 1, 2, \dots \dots, N = 1$$

Sehingga $\psi_{(x)} = U_k(x) \exp(i 2\pi s x / Na)$

↳ untuk pemodelan kromy penny di atas

$$\frac{-\hbar^2}{2m} \frac{\partial^2 \psi}{\partial x^2} + u(x)\psi = E\psi$$

Pada daerah $0 < x < a$ saat $v = 0$ maka:

$$\psi = A e^{ikx} + B e^{-ikx} \text{ dimana } E = \frac{\hbar^2 k^2}{2m}$$

Pada daerah $-b < x < 0$ maka

$$\psi = C e^{Qx} + D e^{-Qx}$$

$$\text{dimana } U_0 = E = \frac{\hbar^2 Q^2}{2m}$$

Dihubungkan solusi $a < x < a+b$ harus dihubungkan pada daerah

$-b < x < 0$

$$\psi(a < x < a + b) = \psi(-b < x < 0) e^{ik(a+b)}$$

Menggunakan syarat batas kontinuitas

$$A + B = C + D$$

$$ik(A - B) = Q(C - D)$$

$$A e^{ika} + B e^{-ika} = (C e^{-Qb} + D e^{Qb}) e^{ik(a+b)}$$

$$ik(A e^{ika} - B e^{-ika}) = Q(C e^{-Qb} - D e^{Qb}) e^{ika}$$

jika $P = \frac{Q^2 b a}{2}$ maka

$$\left(\frac{P}{ka}\right) \sin ka + \cos ka = \cos ka$$

